Transport in interacting disordered nanotubes

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Overview

After introduction:

- Nonlinear magnetotransport in chiral interacting SWNTs (A. De Martino, A. Tsvelik)
- Crossover from Luttinger liquid to Altshuler-Aronov diffusive corrections in MWNTs (C. Mora, A. Altland)

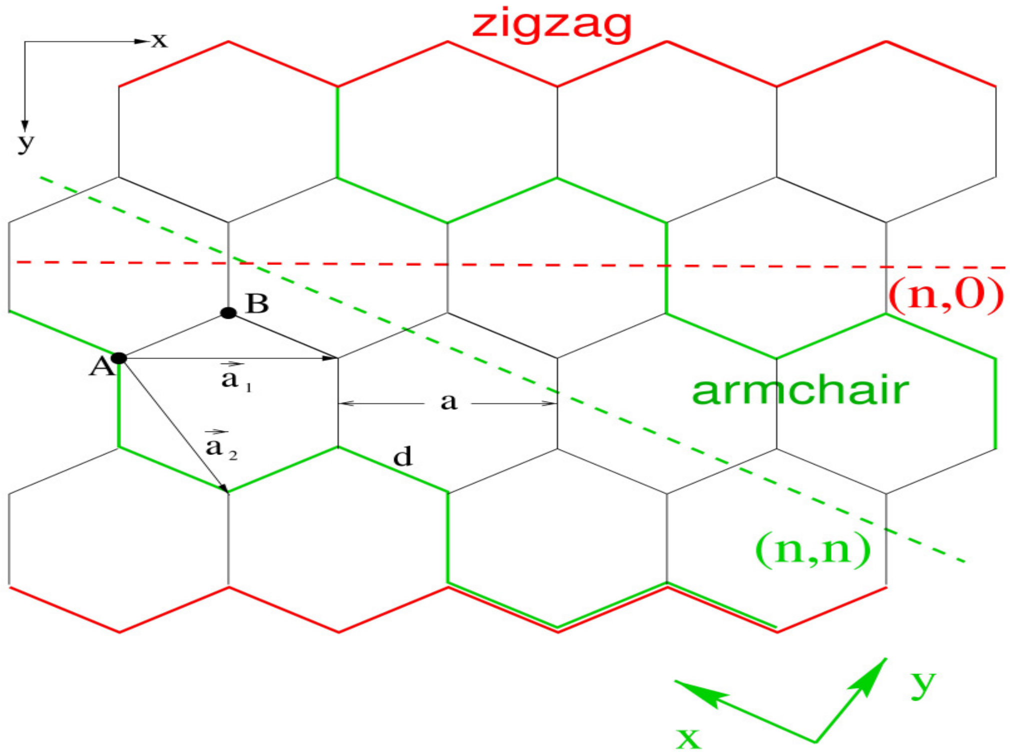
Wrapped 2D graphene sheet

Basis contains two atoms

$$a = \sqrt{3}d\,, d = 0.14nm$$

(n,m) indices: wrapping of sheet onto cylinder

Chiral angle θ: defined with respect to zigzag (n,0) tube

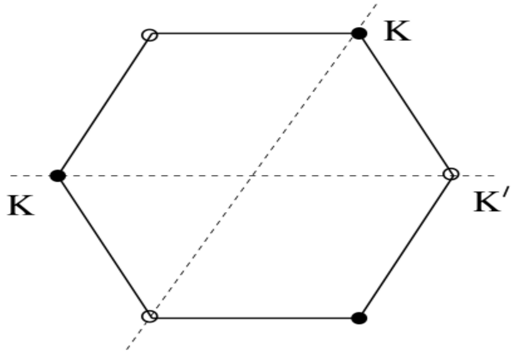


Band structure: Graphene

Exactly two independent corner points K, K´ in first Brillouin zone.

Band structure: valence and conduction bands touch at corner points *(E=0)*, these are the Fermi points in graphene

- Lowest-order k•p scheme: Dirac light cone dispersion
- Deviations at higher energies: trigonal warping



$$E(\vec{q}) = v|\vec{q}|$$

$$\vec{q} = \vec{k} - \vec{K}$$

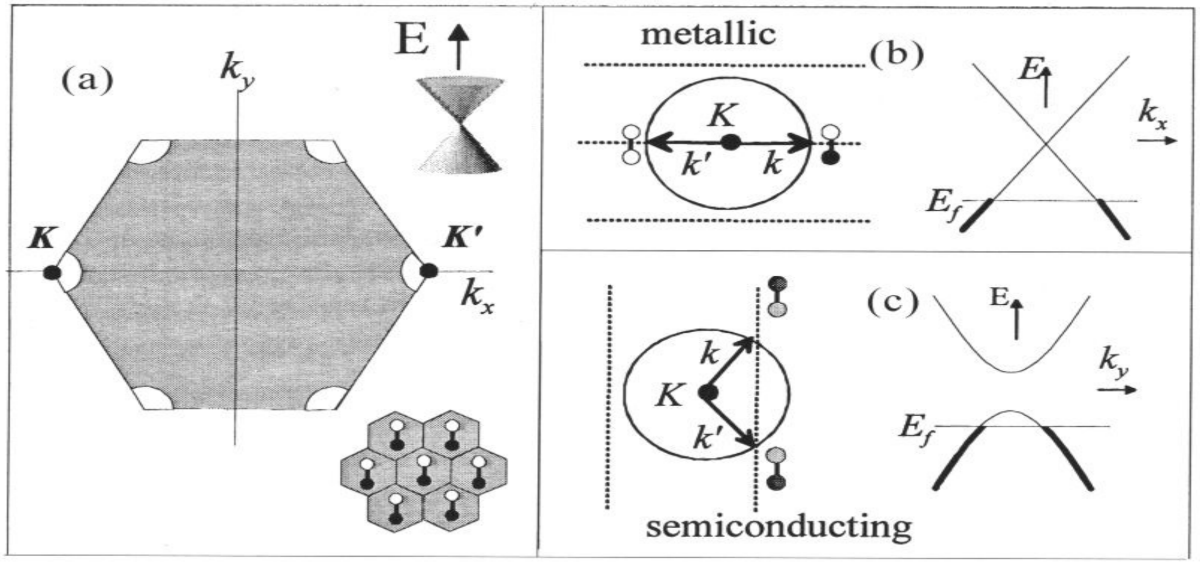
$$v = 8\times 10^{5}\, m/\sec$$

Periodic boundary conditions: SWNTs

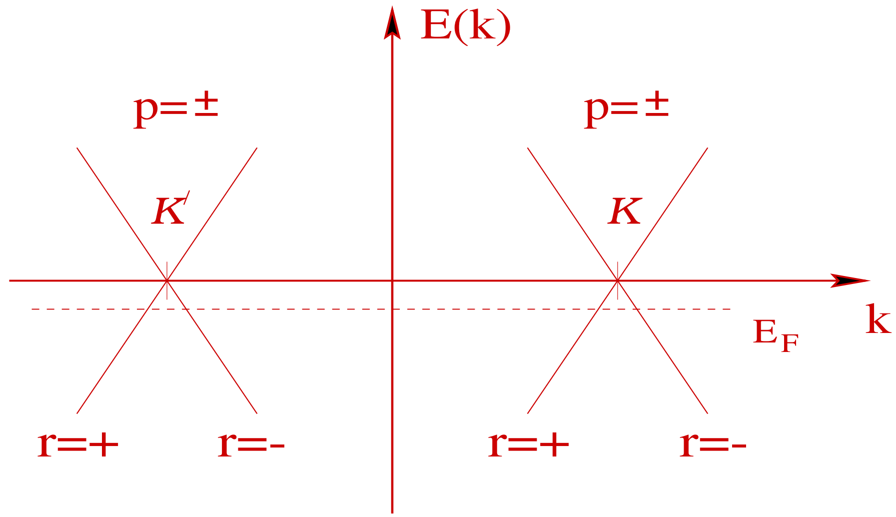
Transverse momentum must be quantized

Nanotube metallic only if K point has allowed transverse momentum

gives necessary condition: *2n+m = 3 x* integer



Metallic SWNTs



- Transverse momentum quantization: keep only $k_{\perp} = 0$
- Ideal 1D quantum wire: 2 spin-degenerate bands
- Low-energy theory: restrict to these 2 bands, but include (long-ranged) Coulomb interactions

Egger & Gogolin, PRL 1997, EPJB 1998
Kane, Balents & Fisher, PRL 1997

Bosonized form

Four bosonic fields, index $a = c+, c-, s+, s-$

Low-energy theory: Luttinger liquid

$$H = \sum_a \frac{v_a}{2} \int dx \left[g_a \Pi_a^2 + g_a^{-1} (\partial_x \varphi_a)^2 \right]$$

$$g_{a \neq c+} \cong 1 \qquad g \equiv g_{c+} \approx 0.2$$

$$v_{c+} = v / g, \; v_{a \neq c+} = v$$

exactly solvable Gaussian model, leads to spin-charge separation and quasi-particles with fractional charge & fractional statistics

Beyond lowest-order k•p scheme?

Dirac cone approximation: chirality drops out

To go beyond, one must include

- Trigonal warping: anisotropic & nonlinear dispersion relation
- Transverse momentum quantization: in parallel magnetic field *B*, including tube curvature

$$k_\perp = eBR^2/2h \pm (a/R)\cos 3\theta$$

- Net effect: R/L movers have different velocity

$$\delta = \frac{v_R - v_L}{v_R + v_L} = \frac{B}{B_0}\sin 6\theta$$

$$B_0 \propto k_F R$$

Nonlinear current-voltage relation

- Linear transport: Onsager-Casimir relation

$$G(B) = G(-B)$$

- Out of equilibrium: odd-in-B part allowed

$$I_e(V,B) = -I_e(V,-B)$$

this contribution is even in voltage!

- Fundamentally interesting because nonzero effect requires combined presence of
  - Electron-electron interactions
  - Chirality (handedness): broken inversion symmetry
  - Magnetic field: broken time reversal symmetry

Sanchez & Büttiker, PRL 2004
Spivak & Zyuzin, PRL 2004

How to include in low energy theory?

- Luttinger liquid theory now comes with R/L moving plasmon velocities, but still exactly solvable Gaussian theory

$$v_{c+,R/L}/v = g^{-1} \pm \delta$$

$$v_{a\neq c+,R/L} = v_{R/L} = v(1 \pm \delta)$$

- Consider long SWNT & good contacts
  - Effect requires at least two impurities
  - Here: 2 impurities separated by distance *d*
  - Nonequilibrium Keldysh approach

Odd-in-B current in a chiral SWNT

De Martino, Egger & Tsvelik, PRL (in press)
cond-mat/0605645

Analytical result:

$$I_e \propto \sin(2k_F d) T_0^{2g-1} e^{-gT_0} \sin\left(\frac{(1-g^2)B}{gB_0} \sin(6\theta) U \right)$$

$$\times \mathrm{Im}\left[e^{iU} \frac{\Gamma(1+g-iU/T_0)}{\Gamma(g)\Gamma(2-iU/T_0)} F\left(g, 1+g-iU/T_0; 2-iU/T_0; e^{-2T_0}\right) \right]$$

with dimensionless temperature/voltage $T_0 = \frac{2\pi k_B T}{\hbar v / gd}, \; U = \frac{|eV|}{\hbar v / gd}$

Requires interactions (g<1) and chirality (sin6θ≠0)

odd in magnetic field *B*, even in bias voltage *V*

changes sign with handedness (enantioselective)

Available experimental results

Measured: $\alpha(T) = \left[\frac{I_e(V,T,B)}{V^2 B} \right]_{V,B \to 0}$

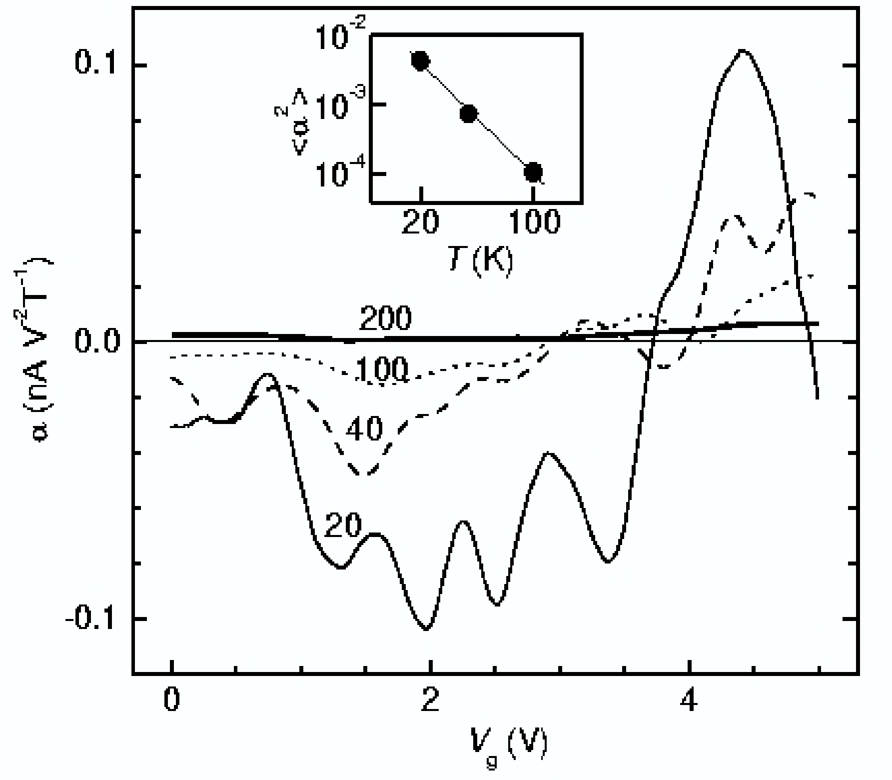
Wei, Cobden et al., PRL 2005

for individual SWNT

- Oscillatory dependence on gate voltage
- Exponentially small at high T, increases when lowering T. Our theory:

$$\alpha(T) \propto T^{(g-1)/2}$$

- Sign does not change with temperature



Oscillations in $I_e(V)$

Zero temperature limit:

$$I_e \propto \sin\left[\frac{(1-g^2)B}{gB_0}\sin(6\theta)U\right]U^{g-1/2}J_{g-1/2}(U)$$

predicts oscillations as function of V with periods:

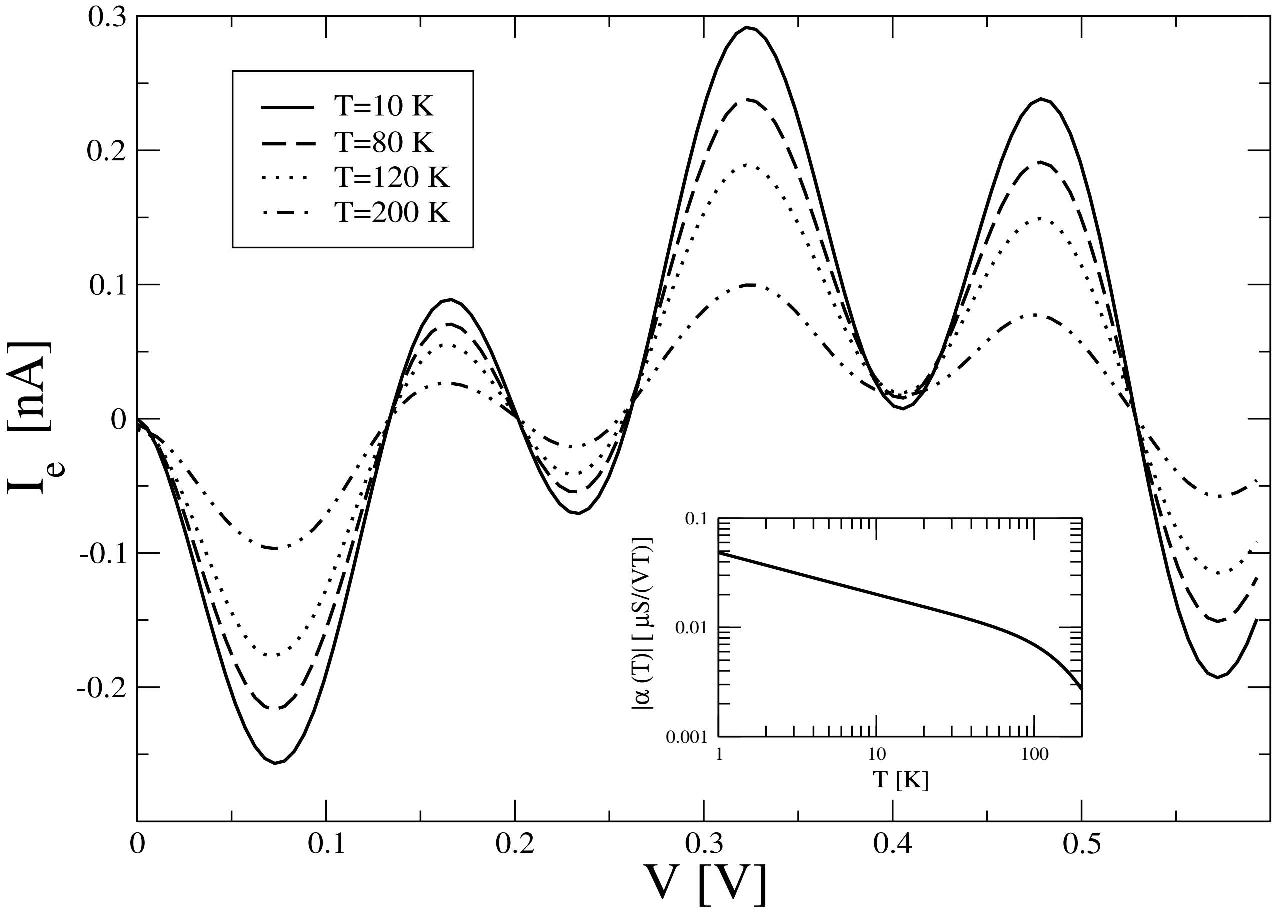
$$\Delta V_1 = \frac{hv}{egd}$$ yields Luttinger parameter

$$\Delta V_2 = \frac{B_0 g \Delta V_1}{B(1-g^2)\sin(6\theta)}$$ yields chirality

Low-voltage limit: Power-law scaling $I_e(V \to 0) \propto |V|^{2g}$

Oscillation with bias voltage

d=20nm
g=0.23
B=16T
(10,4) SWNT



direct observation of interaction physics possible

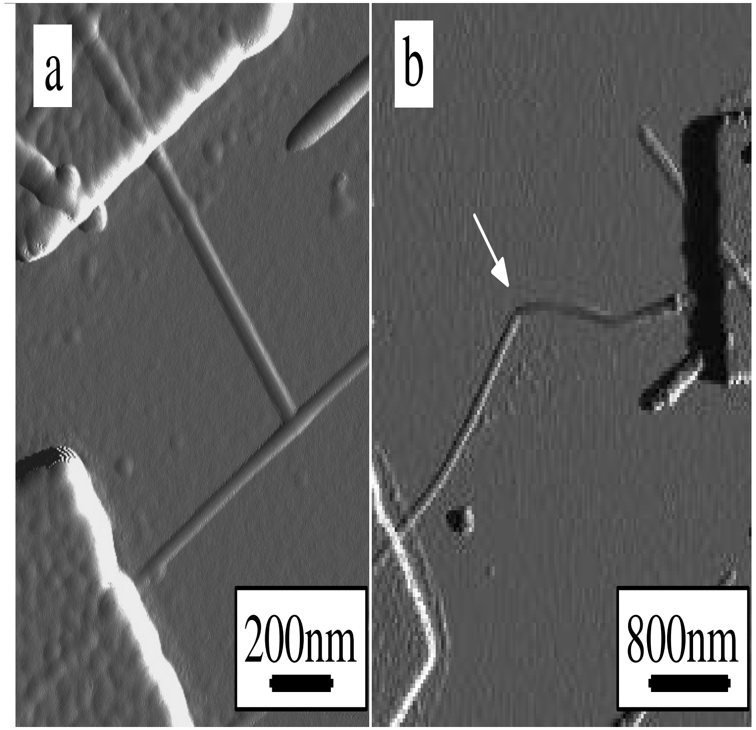
Disorder effects in multichannel wires: Transport in MWNTs

- Electronic transport in MWNTs usually in outermost shell only
- Energy scales one order smaller
- Typically $N \approx 10$ bands due to doping
- Inner shells can also create `disorder´
  - Experiments indicate mean free path $\ell > R$
  - Ballistic behavior for $\omega\tau > 1, \tau = \ell / v$
- Also relevant for long SWNTs

Experiment: TDoS of MWNT

Bachtold et al., PRL 2001

- TDoS observed from conductance through tunnel contact
- Power law zero-bias anomalies
- Scaling properties, similar to Luttinger predictions



TDoS of multi-band Luttinger liquid

Power-law suppressed TDoS reflects orthogonality catastrophe: Electron splinters into true quasiparticles

Geometry dependence

Matveev & Glazman, PRL 1993
Egger, PRL 1999

$$\rho(x,\omega) = \mathrm{Re}\int_0^\infty dt e^{i\omega t}\left\langle \Psi(x,t)\Psi^+(x,0)\right\rangle \propto \omega^\eta$$

$$\eta_{bulk} \equiv \eta = \left(g + 1/g - 2\right)/2N$$

$$\eta_{end} = (1/g - 1)/N > 2\eta$$

Exponents are one order of magnitude too small to explain Basel experiment. Role of disorder?

Interplay of disorder and interaction

Mora, Egger & Altland, cond-mat/0602411

- Coulomb interaction enhanced by disorder
- Expected: crossover from quasiballistic Luttinger liquid at $\omega\tau > 1$ to diffusive/localized phase (e.g. Altshuler-Aronov diffusive anomalies) at $\omega\tau < 1$
- Field theory for multi-channel case and arbitrary disorder strength:

Interacting Nonlinear σ Model

Earlier versions:
Finkel'stein, Z. Phys. B 1983
Kamenev & Andreev, PRB 1999

Bulk TDoS

- Analytical result for $\rho(\omega, T)$ available
- Can be recast in terms of standard *P(E)* Coulomb blockade theory (microscopic derivation)

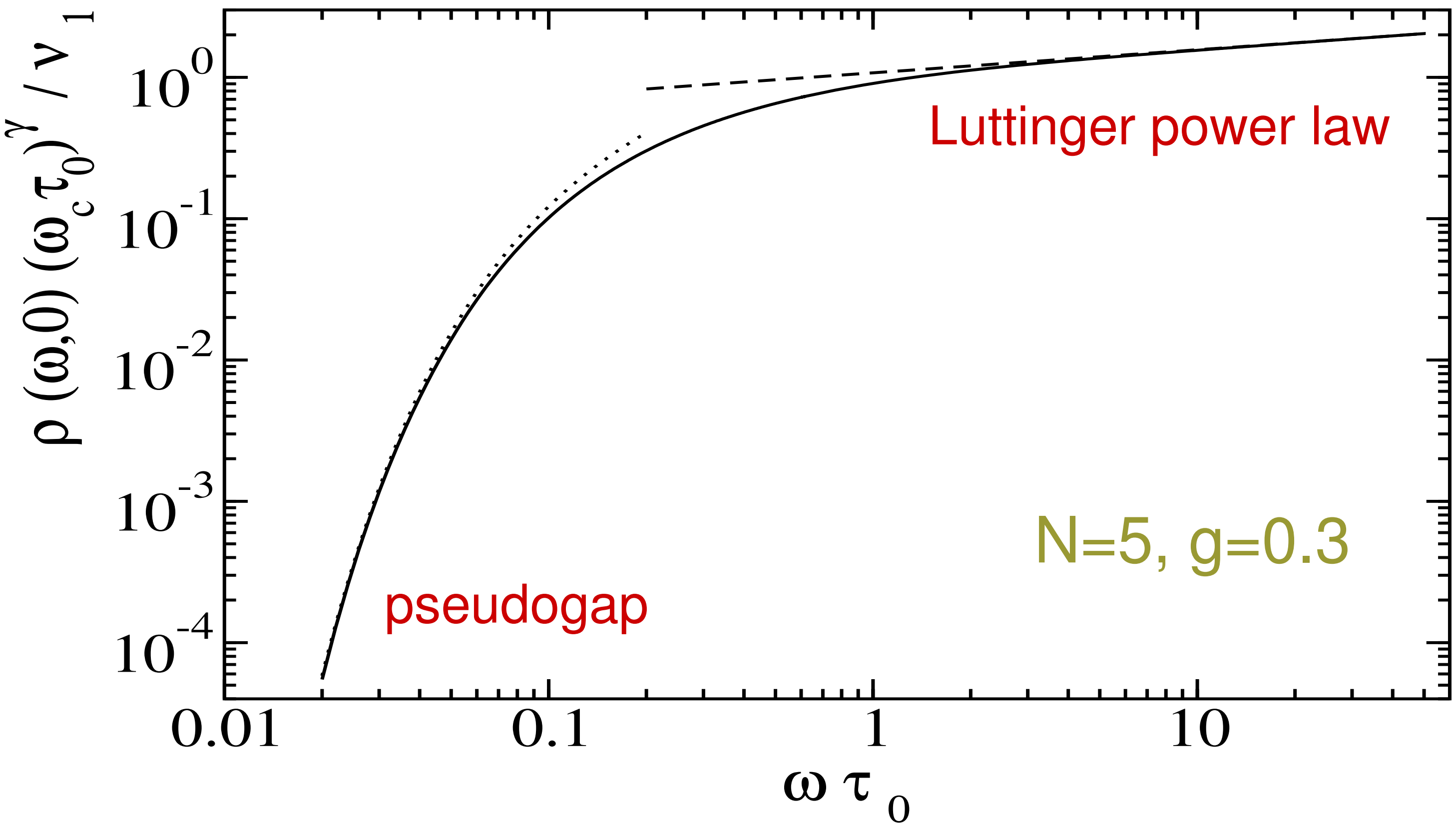
Egger & Gogolin, PRL 2001
Rollbühler & Grabert, PRL 2001

- Zero temperature: describes crossover from
  - Luttinger power law $\rho(\omega\tau > 1) \propto \omega^{\eta}$
  - to pseudogap at low energy:

$$\rho(\omega\tau < 1) \propto \frac{\sqrt{\omega\tau}}{\eta_{end}} \exp\left(-\frac{2\pi\eta_{end}^2}{\omega\tau}\right)$$

Nazarov, JETP 1989
Mishchenko et al., 2001

Bulk TDoS at T=0



Stronger suppression of TDoS due to disorder.
But does not really explain experimental results…

Interaction correction to conductivity

Complete crossover solution from ballistic to diffusive case, to lowest order in interaction:

$$\frac{\sigma(T)}{\sigma_{Drude}} = 1 + \gamma \ln(T\tau') - \gamma \int_0^\infty d\Omega \, \frac{\partial_\Omega \left[\Omega \coth \dfrac{\hbar\Omega}{2k_B T} \right]}{\Omega}$$

$$\times \left((1 + i/\Omega\tau')^{-1/2} - 1 + \frac{i}{2\tau'(1+g)\sqrt{\Omega^2 + i\Omega/\tau'}} \right)$$

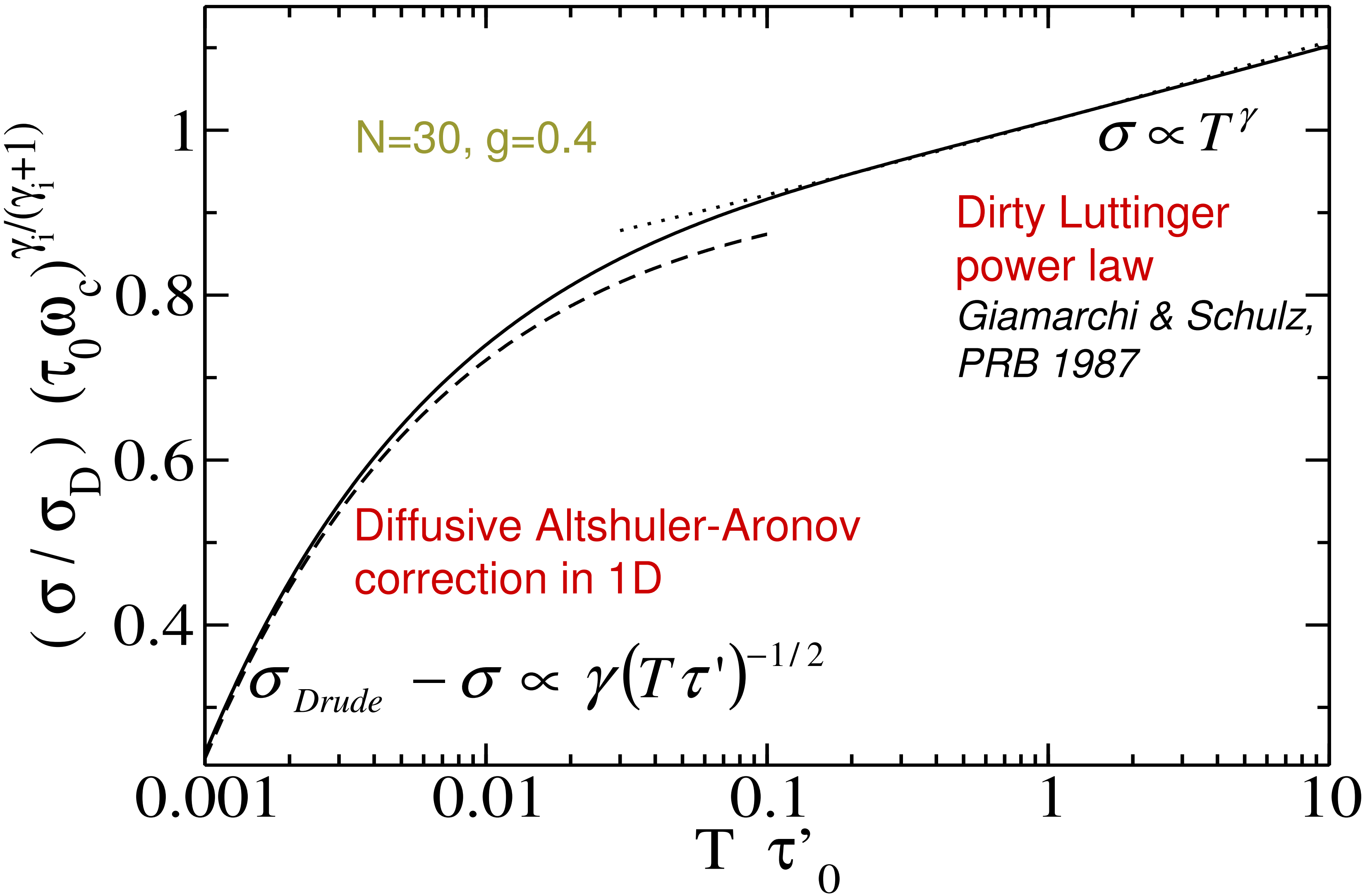
Exponent for weak backscattering by single impurity in a Luttinger liquid

$$\gamma = 2(1-g)/N$$

Renormalized mean free time

$$\tau' = \tau(\omega_c \tau)^{-\gamma/\gamma+1}$$

Conductivity



Conclusions

- Magnetotransport: Linear in B current
  - Only present out of equilibrium (current is even in voltage) and with at least two impurities and with interactions and for chiral tubes!
  - Prediction: Oscillations with bias voltage, power law scaling in T dependence

De Martino, Egger & Tsvelik, cond-mat/0605645

- Crossover from ballistic to diffusive regime: MWNTs
  - Large *N* allows for nonlinear sigma model description: TDoS, conductivity
  - Luttinger liquid and Altshuler-Aronov dominated phase are two sides of the same coin!

Mora, Egger & Altland, cond-mat/0602411